Spontaneous CP Violation at Finite Temperature in the MSSM

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I. Introduction

- 3 requirements to generate BAU : $n_B/s \simeq 10^{-11}$
- (1) baryon number violation
- (2) C and CP violation
- (3) departure from equilibrium
 - EW theory satisfies these if the EWPT is first order.

Electroweak Baryogenesis

 For sufficient CP violation, some extension of the Minimal SM is needed.

for a review,

K.F., Prog.Theor.Phys.96 ('96) 475.

V.A. Rubakov and M.E. Shaposhnikov, hep-ph/9603208.

A.Cohen, et al., Ann.Rev.Nucl.Part.Sci. 33 ('93) 27.

generated BAU (by the charge transport scenario)

$$\frac{n_B}{s} \simeq \mathcal{N} \frac{100}{\pi^2 g_*} \cdot \kappa \alpha_W^4 \cdot \tau T \cdot \frac{F_Y}{v_w T^3}$$
$$\simeq 10^{-3} \times \frac{F_Y}{v_w T^3} \qquad \text{for an optimal case}$$

 F_Y : chiral charge flux $\leftarrow CP$ violation at EW bubble wall $\iff \Delta R$: reflection prob. of chiral fermions



CP violation effective for F_Y at the lowest order

- relative phases of μ, M₂, M₁, A_t in the MSSM chargino, neutralino, stop transport
 [Huet and Nelson, PRD53('96); Aoki, et al. PTP98('97)]
- relative phase θ of the two Higgs doublets

$$\left\langle \Phi_i \right\rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} 0\\ \rho_i e^{i\theta_i} \end{pmatrix}$$

quarks and leptons \Leftarrow Yukawa coupl. $\propto \rho_i e^{i\theta}$ chargino, neutralino, stop mass matrix N.B. $\theta = \theta(x)$ \therefore it cannot be rotated away [Nelson et al. NPB373('92); FKOTT, PRD50('94), PTP95('96)]

CP violating θ at EWPT



a scenario to have large θ near the bubble wall

- = spontaneous CP violation in the transient region
 - + small explicit *CP* violation to resolve degeneracy between *CP* conjugate bubbles

net BAU

$$\frac{n_B}{s} = \frac{\sum_j \left(\frac{n_B}{s}\right)_j N_j}{\sum_j N_j}$$

with

$$N_j = \exp(-4\pi R_C^2 \mathcal{E}_j / T_C)$$

 $\mathcal{E}_j = \text{energy density of the type-} j$ bubble

an example:

$$m_3^2 \longrightarrow m_3^2 e^{-\delta}$$
 with $\delta = 10^{-3}$



II. Transitional CP Violation

Suppose that at
$$T \simeq T_C$$
,
 $V_{\text{eff}}(\rho_i, \theta = \theta_1 - \theta_2)$
 $= \frac{1}{2}m_1^2\rho_1^2 + \frac{1}{2}m_2^2\rho_2^2 - m_3^2\rho_1\rho_2\cos\theta + \frac{\lambda_1}{8}\rho_1^4 + \frac{\lambda_2}{8}\rho_2^4$
 $+ \frac{\lambda_3 + \lambda_4}{4}\rho_1^2\rho_2^2 + \frac{\lambda_5}{4}\rho_1^2\rho_2^2\cos2\theta - \frac{1}{2}(\lambda_6\rho_1^2 + \lambda_7\rho_2^2)\rho_1\rho_2\cos\theta$
 $- [A\rho_1^3 + \rho_1^2\rho_2(B_0 + B_1\cos\theta + B_2\cos2\theta)$
 $+\rho_1\rho_2^2(C_0 + C_1\cos\theta + C_2\cos2\theta) + D\rho_2^3]$
 $= \left[\frac{\lambda_5}{2}\rho_1^2\rho_2^2 - 2(B_2\rho_1^2\rho_2 + C_2\rho_1\rho_2^2)\right]$
 $\times \left[\cos\theta - \frac{2m_3^2 + \lambda_6\rho_1^2 + \lambda_7\rho_2^2 + 2(B_1\rho_1 + C_1\rho_2)}{2\lambda_5\rho_1\rho_2 - 8(B_2\rho_1 + C_2\rho_2)}\right]^2$

 $+\theta$ -independent terms

where all the parameters are real

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conditions for the transitional CP violation for a given (ρ_1, ρ_2)

$$\begin{split} F(\rho_1,\rho_2) &\equiv \frac{\lambda_5}{2} \rho_1^2 \rho_2^2 - 2(B_2 \rho_1^2 \rho_2 + C_2 \rho_1 \rho_2^2) > 0, \\ -1 < G(\rho_1,\rho_2) &\equiv \frac{2m_3^2 + \lambda_6 \rho_1^2 + \lambda_7 \rho_2^2 + 2(B_1 \rho_1 + C_1 \rho_2)}{2\lambda_5 \rho_1 \rho_2 - 8(B_2 \rho_1 + C_2 \rho_2)} < 1 \\ \text{At } T \simeq T_C, \text{ around the EW bubble wall} \end{split}$$

$$(\rho_1, \rho_2)$$
 : $(0, 0) \longrightarrow (v_C \cos \beta_C, v_C \sin \beta_C)$

There may be a chance to satisfy the conditions in the transient region.

III. Effective parameters of the MSSM

MSSM at the tree level:

$$\lambda_1 = \lambda_2 = \frac{1}{4}(g_2^2 + g_1^2), \quad \lambda_3 = \frac{1}{4}(g_2^2 - g_1^2), \quad \lambda_4 = -\frac{1}{2}g_2^2,$$

$$\lambda_5 = \lambda_6 = \lambda_7 = 0,$$

the relevant parameters in F and G are induced by radiative and finite-temperature corrections

effective potential at the one-loop level

$$V_{\text{eff}} = V_0 + V_1(\rho_i, \boldsymbol{\theta}) + \bar{V}_1(\rho_i, \boldsymbol{\theta}; T),$$

where

with

$$\varphi_d = \frac{1}{\sqrt{2}} \begin{pmatrix} \rho_1 \\ 0 \end{pmatrix} = \frac{1}{\sqrt{2}} \begin{pmatrix} v_1 \\ 0 \end{pmatrix},$$
$$\varphi_u = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ \rho_2 e^{i\theta} \end{pmatrix} = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ v_2 + iv_3 \end{pmatrix}$$

We consider the contributions from the charginos and neutralinos (χ) , stops (\tilde{t}) and charged Higgs bosons (ϕ^{\pm}) .

relevant effective parameters:

$$\begin{aligned} \left(m_3^2\right)_{\text{eff}} &= -\frac{\partial^2 V_{\text{eff}}}{\partial v_1 \partial v_2} \Big|_0 = m_3^2 + \Delta_{\chi} m_3^2 + \Delta_{\tilde{t}} m_3^2 + \Delta_{\phi^{\pm}} m_3^2, \\ \lambda_5 &= \frac{1}{2} \left(\frac{\partial^4 V_{\text{eff}}}{\partial v_1^2 \partial v_2^2} \Big|_0 - \frac{\partial^4 V_{\text{eff}}}{\partial v_1^2 \partial v_3^2} \Big|_0 \right), \\ \lambda_6 &= -\frac{1}{3} \left. \frac{\partial^4 V_{\text{eff}}}{\partial v_1^3 \partial v_2} \right|_0 = \Delta_{\chi} \lambda_6 + \Delta_{\tilde{t}} \lambda_6 + \Delta_{\phi^{\pm}} \lambda_6, \\ \lambda_7 &= -\frac{1}{3} \left. \frac{\partial^4 V_{\text{eff}}}{\partial v_1 \partial v_2^3} \right|_0 = \Delta_{\chi} \lambda_7 + \Delta_{\tilde{t}} \lambda_7 + \Delta_{\phi^{\pm}} \lambda_7 \end{aligned}$$

 $\rho^3\text{-terms}$ in $V_{\rm eff}$ \Longleftarrow zero modes of the bosons at $T\neq 0$ $W^\pm,~Z,~{\rm Higgs,~squarks,~sleptons}$

$$\begin{array}{l} \theta \text{-dependent } \rho^3 \text{-terms } (B_i \text{ and } C_i) \\ \longleftarrow m_{\tilde{t}} \simeq 0 \\ M_{\tilde{t}}^2 = \begin{pmatrix} m_{\tilde{q}}^2 + m_t^2 - \frac{3g_2^2 - g_1^2}{8}(\rho_1^2 - \rho_2^2) & \frac{y_t}{\sqrt{2}}(\mu\rho_1 + A_t\rho_2 e^{-i\theta}) \\ \frac{y_t}{\sqrt{2}}(\mu\rho_1 + A_t\rho_2 e^{i\theta}) & m_{\tilde{t}}^2 + m_t^2 + \frac{g_1^2}{6}(\rho_1^2 - \rho_2^2) \end{pmatrix} \end{array}$$

and

$$\bar{V}_{\tilde{t}}(\rho_i, \theta; T) = 3 \frac{T^4}{2\pi^2} \left[2I_B(a_+^2) + 2I_B(a_-^2) \right],$$

with

$$\begin{split} I_B(a_-^2) &= \int_0^\infty dx \, x^2 \log \left(1 - e^{-\sqrt{x^2 + a_-^2}} \right) \\ &= -\frac{\pi^4}{45} + \frac{\pi^2}{12} a_-^2 - \frac{\pi}{6} a_-^3 + \lambda_- a_-^4 + \dots \\ a_-^2 &\equiv m_-^2/T^2 \qquad ; \qquad m_-^2(\rho) = \text{smaller eigenvalue of } M_{\tilde{t}}^2 \end{split}$$

• charginos and neutralinos

$$M_{\chi^{\pm}} = \begin{pmatrix} M_2 & -\frac{ig_2}{\sqrt{2}}\rho_2 e^{-i\theta} \\ -\frac{ig_2}{\sqrt{2}}\rho_1 & -\mu \end{pmatrix},$$

$$M_{\chi^0} = \begin{pmatrix} M_2 & 0 & -\frac{i}{2}g_2\rho_1 & \frac{i}{2}g_2\rho_2 e^{-i\theta} \\ 0 & M_1 & \frac{i}{2}g_1\rho_1 & -\frac{i}{2}g_1\rho_2 e^{-i\theta} \\ -\frac{i}{2}g_2\rho_1 & \frac{i}{2}g_1\rho_1 & 0 & \mu \\ \frac{i}{2}g_2\rho_2 e^{-i\theta} & -\frac{i}{2}g_1\rho_2 e^{-i\theta} & \mu & 0 \end{pmatrix}$$

chargino and neutralino contributions ($M_2 = M_1$)

$$\begin{split} \Delta_{\chi} m_3^2 &= -2g_2^2 \left(1 + \frac{1}{\cos^2 \theta_W} \right) \mu M_2 L(M_2, \mu) \\ &+ \frac{g_2^2}{\pi^2} \left(1 + \frac{1}{\cos^2 \theta_W} \right) \mu M_2 f_2^{(+)} \left(\frac{M_2}{T}, \frac{\mu}{T} \right), \\ \Delta_{\chi} \lambda_5 &= \frac{g_2^4}{8\pi^2} \left(1 + \frac{2}{\cos^4 \theta_W} \right) K \left(\frac{M_2^2}{\mu^2} \right) \\ &- \frac{g_2^4}{\pi^2 T^4} \left(1 + \frac{2}{\cos^4 \theta_W} \right) \mu^2 M_2^2 f_4^{(+)} \left(\frac{M_2}{T}, \frac{\mu}{T} \right), \\ \Delta_{\chi} \lambda_6 &= -\frac{g_2^4}{8\pi^2} \left(1 + \frac{2}{\cos^4 \theta_W} \right) \frac{\mu}{M_2} \left[-H \left(\frac{M_2^2}{\mu^2} \right) + K \left(\frac{M_2^2}{\mu^2} \right) \right] \\ &+ \frac{g_2^4}{\pi^2} \left(1 + \frac{2}{\cos^4 \theta_W} \right) \left[\frac{\mu M_2}{T^2} f_3^{(+)} \left(\frac{M_2}{T}, \frac{\mu}{T} \right) \right] \\ &+ \frac{\mu^3 M_2}{T^4} f_4^{(+)} \left(\frac{M_2}{T}, \frac{\mu}{T} \right) \right] = \Delta_{\chi} \lambda_7 \end{split}$$

where

$$L(m_1, m_2) = \frac{1}{16\pi^2} \left[1 - \frac{m_1^2}{m_1^2 - m_2^2} \log \frac{m_1^2}{M_{\text{ren}}^2} + \frac{m_2^2}{m_1^2 - m_2^2} \log \frac{m_2^2}{M_{\text{ren}}^2} \right]$$

> 0,
$$H(\alpha) = \frac{\alpha}{\alpha - 1} \left(\frac{1}{\alpha - 1} \log \alpha - 1 \right) < 0,$$
$$K(\alpha) = \frac{\alpha}{(\alpha - 1)^2} \left(\frac{\alpha + 1}{\alpha - 1} \log \alpha - 2 \right) > 0,$$

and

$$\begin{split} f_2^{(\mp)}(a,b) &= -\frac{1}{a^2 - b^2} \int_0^\infty dx (\frac{x^2}{\sqrt{x^2 + a^2}} \frac{1}{e^{\sqrt{x^2 + a^2}}} + 1) \\ &- \frac{x^2}{\sqrt{x^2 + b^2}} \frac{1}{e^{\sqrt{x^2 + b^2}}} + 1) > 0, \\ f_3^{(\mp)}(a,b) &= \frac{1}{2(a^2 - b^2)} \int_0^\infty \frac{dx}{\sqrt{x^2 + a^2}} \frac{1}{e^{\sqrt{x^2 + a^2}}} + 1 \\ &- \frac{1}{a^2 - b^2} f_2^{(\mp)}(a,b) < 0, \\ f_4^{(\mp)}(a,b) &= \frac{1}{2(a^2 - b^2)^2} \int_0^\infty dx (\frac{1}{\sqrt{x^2 + a^2}} \frac{1}{e^{\sqrt{x^2 + a^2}}} + 1) \\ &+ \frac{1}{\sqrt{x^2 + b^2}} \frac{1}{e^{\sqrt{x^2 + b^2}}} + 1 \\ &- \frac{2}{(a^2 - b^2)^2} f_2^{(\mp)}(a,b) > 0, \end{split}$$

 \bullet stops with $m_{\tilde{t}}\simeq 0$

$$\begin{split} \Delta_{\tilde{t}}m_{3}^{2} &= N_{c}y_{t}^{2}\mu A_{t} L(m_{\tilde{q}},0) + \frac{3T^{2}}{\pi^{2}} \frac{y_{t}^{2}\mu A_{t}}{m_{\tilde{q}}^{2}} \Big[I_{B}^{\prime}(a_{\tilde{q}}^{2}) - \frac{\pi^{2}}{12} \Big] ,\\ \Delta_{\tilde{t}}\lambda_{5} &= -\frac{N_{c}y_{t}^{4}}{16\pi^{2}m_{\tilde{q}}^{2}M_{IR}^{2}} K\left(\frac{m_{\tilde{q}}^{2}}{M_{IR}^{2}}\right) \\ &+ \frac{N_{c}y_{t}^{4}\mu^{2}A_{t}^{2}}{\pi^{2}(m_{\tilde{q}}^{2})^{2}} \Big[\frac{2T^{2}}{m_{\tilde{q}}^{2}} \left(-I_{B}^{\prime}(a_{\tilde{q}}^{2}) + \frac{\pi^{2}}{12} \right) + I_{B}^{\prime\prime}(a_{\tilde{q}}^{2}) + 2\lambda_{-} \Big] \\ \Delta_{\tilde{t}}\lambda_{6} &= \frac{N_{c}y_{t}^{2}\mu A_{t}}{16\pi^{2}m_{\tilde{q}}^{2}} \Big[\frac{1}{4} \left(\frac{g_{1}^{2}}{3} - g_{2}^{2} \right) - \frac{g_{1}^{2}m_{\tilde{q}}^{2}}{3M_{IR}^{2}} H\left(\frac{M_{IR}^{2}}{m_{\tilde{q}}^{2}} \right) \\ &+ \frac{y_{t}^{2}\mu^{2}}{M_{IR}^{2}} K\left(\frac{m_{\tilde{q}}}{M_{IR}^{2}} \right) \Big] \\ &+ \frac{N_{c}y_{t}^{2}\mu A_{t}}{\pi^{2}m_{\tilde{q}}^{2}} \Big\{ \frac{2T^{2}}{m_{\tilde{q}}^{2}} \left(\frac{y_{t}^{2}\mu^{2}}{m_{\tilde{q}}^{2}} + \left(-\frac{5}{3}g_{1}^{2} + g_{2}^{2} \right) \right) \Big[I_{B}^{\prime}(a_{\tilde{q}}^{2}) - \frac{\pi^{2}}{12} \Big] \\ &- \Big(\frac{y_{t}^{2}\mu^{2}}{m_{\tilde{q}}^{2}} + \frac{3g_{2}^{2} - g_{1}^{2}}{12} \Big) I_{B}^{\prime\prime}(a_{\tilde{q}}^{2}) + 2 \Big(\frac{g_{1}^{2}}{3} - \frac{y_{t}^{2}\mu^{2}}{m_{\tilde{q}}^{2}} \Big) \lambda_{-} \Big\} , \\ \Delta_{\tilde{t}}\lambda_{7} &= \frac{N_{c}y_{t}^{2}\mu A_{t}}{16\pi^{2}m_{\tilde{q}}^{2}} \Big[- \Big(y_{t}^{2} + \frac{1}{4} \Big(\frac{g_{1}^{2}}{3} - g_{2}^{2} \Big) \Big) \\ &- \Big(y_{t}^{2} - \frac{g_{1}^{2}}{3} \Big) \frac{m_{\tilde{q}}^{2}}{M_{IR}^{2}} H\left(\frac{M_{IR}}{m_{\tilde{q}}^{2}} \right) + \frac{y_{t}^{2}A_{t}^{2}}{M_{IR}^{2}} K\left(\frac{m_{\tilde{q}}^{2}}{M_{IR}^{2}} \right) \Big] \\ &+ \frac{N_{c}y_{t}^{2}\mu A_{t}}{\pi^{2}m_{\tilde{q}}^{2}} \Big\{ \frac{2T^{2}}{m_{\tilde{q}}^{2}} \left(\frac{y_{t}^{2}\mu^{2}}{m_{\tilde{q}}^{2}} - \left(-\frac{5}{3}g_{1}^{2} + g_{2}^{2} \right) \Big) \Big[I_{B}^{\prime}(a_{\tilde{q}}^{2}) - \frac{\pi^{2}}{12} \Big] \\ &- \Big(y_{t}^{2} + \frac{y_{t}^{2}\mu^{2}}{m_{\tilde{q}}^{2}} \Big] \Big\}$$

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$$+2\left(\ y_t^2 - rac{g_1^2}{3} - rac{y_t^2 \mu^2}{m_{ ilde{q}}^2}
ight) \ \lambda_-
ight\}$$

and for $A_t/m_{\tilde{q}}, \mu/m_{\tilde{q}}, g_1/y_t \ll 1$,

$$C_{1} \simeq \frac{T}{4\sqrt{2}\pi} |y_{t}|^{3} \frac{3\mu A_{t}}{m_{\tilde{q}}^{2}} \left[-1 + \frac{1}{2} \left(\frac{A_{t}}{m_{\tilde{q}}} \right)^{2} + \frac{1}{6} \left(\frac{g_{1}}{y_{t}} \right)^{2} \right] ,$$

$$B_{2} \simeq \frac{T}{4\sqrt{2}\pi} |y_{t}|^{3} \frac{3}{2} \left(\frac{\mu A_{t}}{m_{\tilde{q}}^{2}} \right)^{2} .$$

• charged Higgs bosons

$$\begin{split} \Delta_{\phi^{\pm}} m_3^2 &= \frac{1}{2} g_2^2 m_3^2 L(\mu_1, \mu_2) + \frac{1}{4\pi^2} g_2^2 m_3^2 f_2^{(-)} \Big(\frac{\mu_1}{T}, \frac{\mu_2}{T} \Big) \ , \\ \Delta_{\phi^{\pm}} \lambda_5 &= -\frac{g_2^4}{64\pi^2} \frac{m_3^4}{\mu_1^2 \mu_2^2} K \Big(\frac{\mu_1^2}{\mu_2^2} \Big) - \frac{1}{8\pi^2 T^4} g_2^4 m_3^4 f_4^{(-)} \Big(\frac{\mu_1}{T}, \frac{\mu_2}{T} \Big) \ , \\ \Delta_{\phi^{\pm}} \lambda_6 &= \frac{g_2^4}{64\pi^2} \frac{m_3^2}{\mu_1^2} \Big\{ -H \Big(\frac{\mu_1^2}{\mu_2^2} \Big) + \Big[1 - \frac{m_1^2}{2\mu^2 \cos^2 \theta_W} \\ &- \Big(1 - \frac{1}{2\cos^2 \theta_W} \Big) \frac{m_2^2}{\mu_2^2} \Big] K \Big(\frac{\mu_1^2}{\mu_2^2} \Big) \Big\} \\ &+ \frac{g_2^4 m_3^2}{8\pi^2 T^2} \Big[f_3^{(-)} \Big(\frac{\mu_1}{T}, \frac{\mu_2}{T} \Big) + \Big(\frac{\mu_2^2}{T^2} - \frac{m_1^2}{2T^2 \cos^2 \theta_W} \\ &- \Big(1 - \frac{1}{2\cos^2 \theta_W} \Big) \frac{m_2^2}{T^2} \Big) f_4^{(-)} \Big(\frac{\mu_1}{T}, \frac{\mu_2}{T} \Big) \Big] \ , \\ \Delta_{\phi^{\pm}} \lambda_7 &= \frac{g_2^4}{64\pi^2} \frac{m_3^2}{\mu_1^2} \Big\{ -H \Big(\frac{\mu_1^2}{\mu_2^2} \Big) + \Big[1 - \Big(1 - \frac{1}{2\cos^2 \theta_W} \Big) \frac{m_1^2}{\mu_2^2} \\ &- \frac{m_2^2}{2\mu_2^2 \cos^2 \theta_W} \Big] K \Big(\frac{\mu_1^2}{\mu_2^2} \Big) \Big\} \end{split}$$

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$$\begin{aligned} + \frac{g_2^4 m_3^2}{8\pi^2 T^2} \Big[f_3^{(-)} \Big(\frac{\mu_1}{T}, \frac{\mu_2}{T} \Big) &+ (\frac{\mu_2^2}{T^2} - \Big(1 - \frac{1}{2\cos^2\theta_W} \Big) \frac{m_1^2}{T^2} \\ &- \frac{m_2^2}{2T^2\cos^2\theta_W} \Big) f_4^{(-)} \Big(\frac{\mu_1}{T}, \frac{\mu_2}{T} \Big) \Big] , \end{aligned}$$
$$\mu_{1,2}^2 = \frac{\bar{m}_1^2 + \bar{m}_2^2 \pm \sqrt{(\bar{m}_1^2 - \bar{m}_2^2)^2 + 4m_3^4}}{2} \end{aligned}$$

and

where

$$\bar{m}_1^2 = m_1^2 + \frac{1}{16\pi^2} (3g_2^2 + g_1^2)T^2,$$

$$\bar{m}_2^2 = m_2^2 + \frac{1}{16\pi^2} (3g_2^2 + g_1^2 + 4y_t^2)T^2.$$

III. Numerical Results



 $B_2/T = 7.368276 \times 10^{-6}, \qquad C_1/T = -2.313638 \times 10^{-3}$ $\implies F(\rho_1, \rho_2) > 0 \text{ for } \rho_2 > 5.104 \text{ at } T = 100$

region in which $|G(\rho_1, \rho_2)| < 1$ is satisfied:



(B) $\mu A_t < 0$

m_{3}^{2}	A_t	M_2	μ	$m_{ ilde q}$
2200	10	300	-300	400





IV. Discussions

The conditions for the transitional CP violation:

$$\begin{split} F(\rho_1,\rho_2) > 0 \quad \text{and} \quad |G(\rho_1,\rho_2)| < 1 \\ \text{are satisfied at } T \simeq T_C \text{ if} \\ 1. \ |\mu| \simeq |M_2| \ [\text{for } \Delta_\chi \lambda_5 > 0] \\ 2. \ \mu M_2 < 0 \quad [\text{to decrease } (m_3^2)_{\text{eff}} \text{ by the } \chi\text{-contributions}] \\ 3. \ m_3^2 \text{ is properly tuned at the tree-level} \\ \quad [\text{for } (m_3^2)_{\text{eff}} \text{ to become as small as } \lambda_{6,7}\rho^2] \end{split}$$

This mechanism of CP violation

- is free from the constraints on CP violation at T = 0,
- can generate sufficient BAU,
- is not bothered by the light scalar. At T = 0, spontaneous CP viol. does not occur \implies no pseudo-Goldstone boson $m_A^{\text{tree}} = \sqrt{m_3^2(\tan\beta_0 + \cot\beta_0)} = 107\text{GeV}$ for case (B)

N.B.

- CP violating minimum of $V_{\rm eff}$ in the transient region need not be the global minimum.
- We still need to know the global structure of V_{eff} to determine T_C and the profile of the EW bubble wall.

··· numerical studies in progress