
Signature of the MSSM with ν_{RS} in a long baseline experiment

*Results are shown in part in the proceedings for the conference NuFACT04,
[arXiv:hep-ph/0410408]*

*Paper is in progress [arXiv:hep-ph/0502***]*

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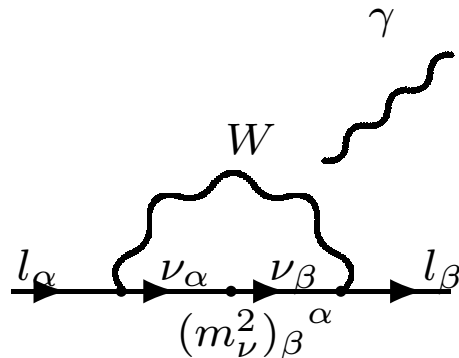
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 - cLFV: lepton flavor violation in charged lepton sector
 - nLFV: lepton flavor violation with neutrinos
- Model independent analysis
 - Feasibility study to detect nLFV signals
based on TO Sato Yamashita, *Phys. Rev.* **D65** (2002) 093015.
- In the MSSM with ν_{RS} ,
not only cLFV but also nLFV are induced.
 - Correlation between cLFV and nLFV
 - Size of nLFV couplings. — Detectable or not?
- Summary

Introduction

Introduction: LFV

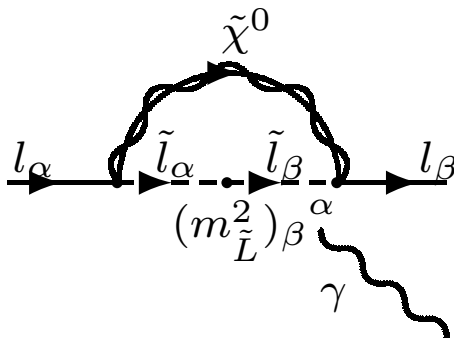
- In the SM, there is no LFV. In the SM with m_ν , it is small.



$$\text{Br}(l_\alpha \rightarrow l_\beta \gamma) \simeq \frac{3\alpha}{32\pi} \left| U_{\mu i} \frac{m_{\nu i}^2}{m_W^2} U_{ie}^\dagger \right|^2$$

Cheng Li (1977), Petcov (1977),
Marciano Sandra (1977), Shrock Lee (1977)

- In the MSSM with ν_{RS} , it can be large
 - The large mixings of neutrinos may imply the large cLFV ...



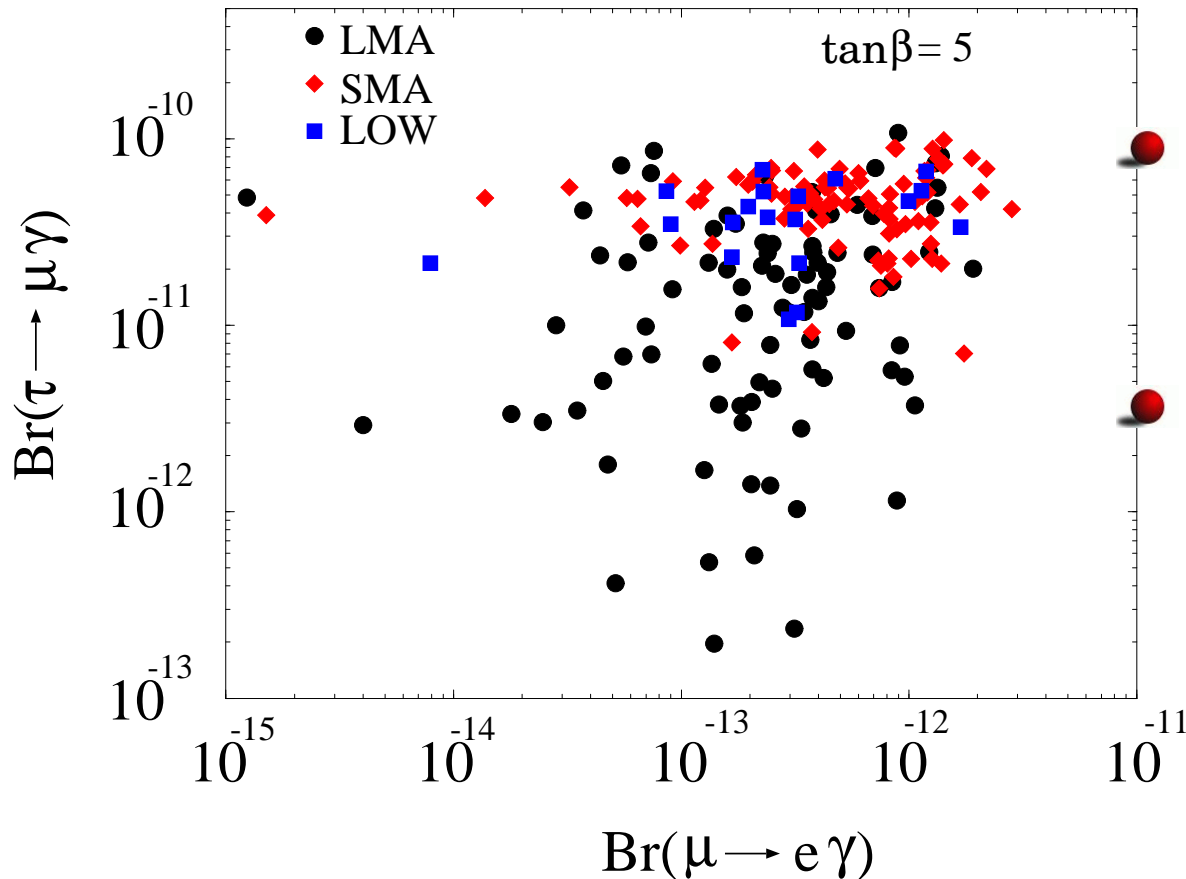
$$\left(m_{\tilde{L}}^2 \right)_\beta^\alpha \simeq -\frac{(6 + a_0^2)m_0^2}{16\pi^2} \left(Y_\nu^\dagger Y_\nu \right)_\beta^\alpha \ln \frac{M_X}{M_R},$$

$$\text{Br}(l_\alpha \rightarrow l_\beta \gamma) \simeq \frac{\alpha^3}{G_F^2} \frac{\left| \left(m_{\tilde{L}}^2 \right)_\beta^\alpha \right|^2}{m_{\text{SUSY}}^8} \tan^2 \beta.$$

Borzumati Masiero (1986),
Hisano Moroi Tobe Yamaguchi (1995)

LFV in charged lepton sector (cLFV)

$M_2 = 150 \text{ GeV}$ $m_{\tilde{e}_L} = 300 \text{ GeV}$ $a_0 = 0$ $\mu > 0$



● One example numerical study

Sato Tobe (2000)

● The future experiments

$\text{Br}(\mu \rightarrow e\gamma) > 10^{-14}$,

$\text{Br}(\tau \rightarrow \mu\gamma) > 10^{-8}$,

$\text{Br}(\mu \rightarrow e \text{ conv.}) > 10^{-18}$.

- The search for the cLFV is promising experiments.
 - However, we here consider an alternative process.

LFV interaction with neutrinos (nLFV)

- We here discuss a process with the neutrino flavor violation, such as

$$\mu^- \rightarrow \nu_\tau e^- \bar{\nu}_e, \quad \nu_\alpha e^- \rightarrow \nu_\beta e^-, \quad \nu_\alpha d \rightarrow \ell_\beta u.$$

These processes affect the neutrino oscillation experiments.

Gonzalez-Garcia Grossman Gusso Nir (2001)

Gago Guzzo Nunokawa Teves Zukanovich Funchal (2001)

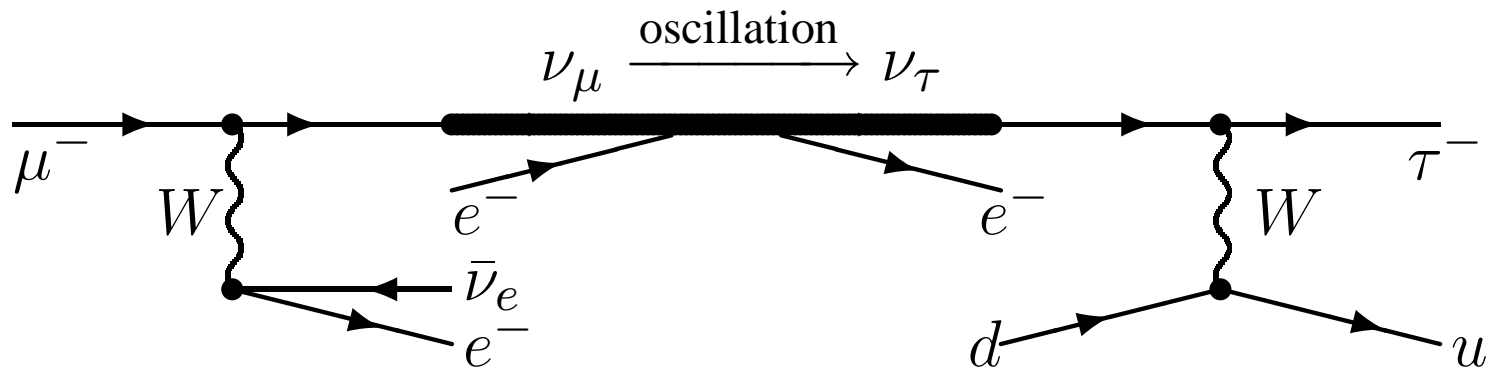
Huber Schwetz Valle (2002)

Fogli Lisi Mirizzi Montanino (2002)

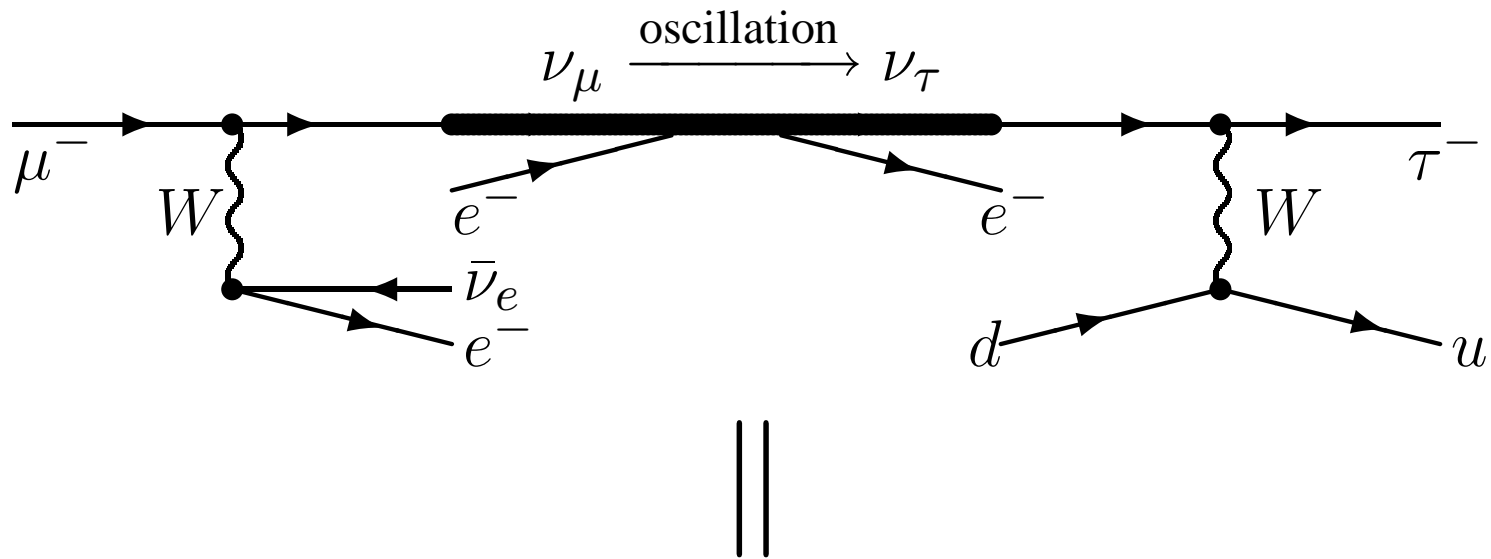
- Can we detect these effects in oscillation experiments?
 - Yes, we can, but it depends on the type and size of nLFV interaction.
- In the MSSM+ ν_R , what is the typical size of nLFV couplings?
 - We show the correlation between nLFV and cLFV.

Model independent approach

Standard Oscillation $\nu_\mu \rightarrow \nu_\tau$

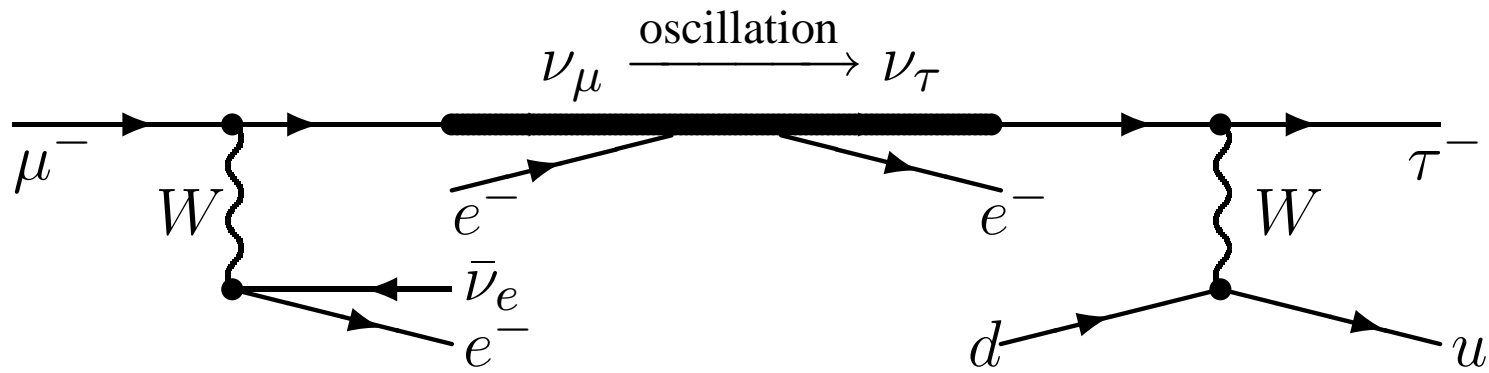


Standard Oscillation $\nu_\mu \rightarrow \nu_\tau$



$$\mathcal{A}(\mu \rightarrow \nu_\mu \bar{\nu}_e e) \times \mathcal{A}(\nu_\mu \xrightarrow{\text{osc}} \nu_\tau) \times \mathcal{A}(\nu_\tau d \rightarrow \tau u)$$

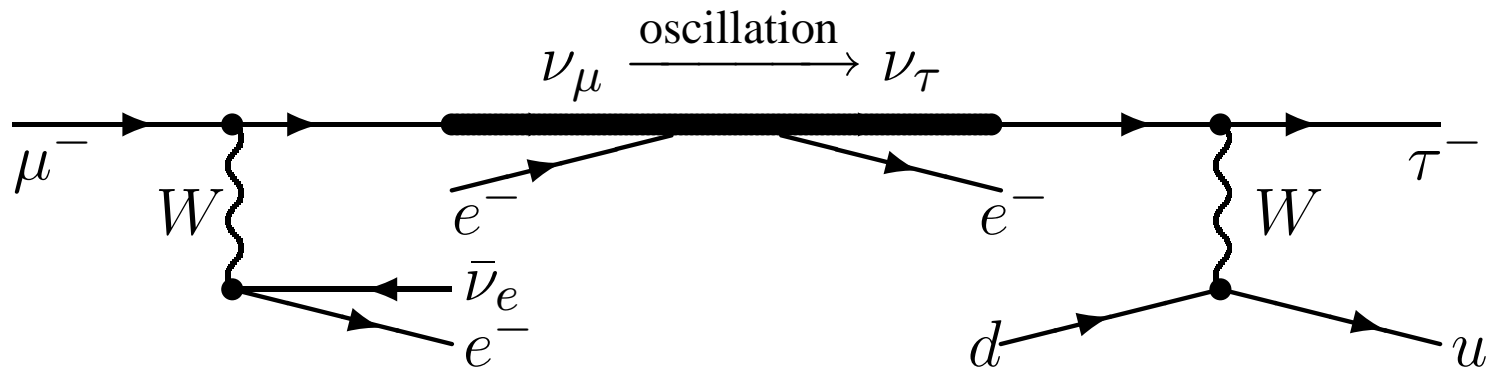
Standard Oscillation $\nu_\mu \rightarrow \nu_\tau$



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$$\left| \mathcal{A}(\mu \rightarrow \nu_\mu \bar{\nu}_e e) \times \mathcal{A}(\nu_\mu \xrightarrow{\text{osc}} \nu_\tau) \times \mathcal{A}(\nu_\tau d \rightarrow \tau u) \right|^2$$

Standard Oscillation $\nu_\mu \rightarrow \nu_\tau$

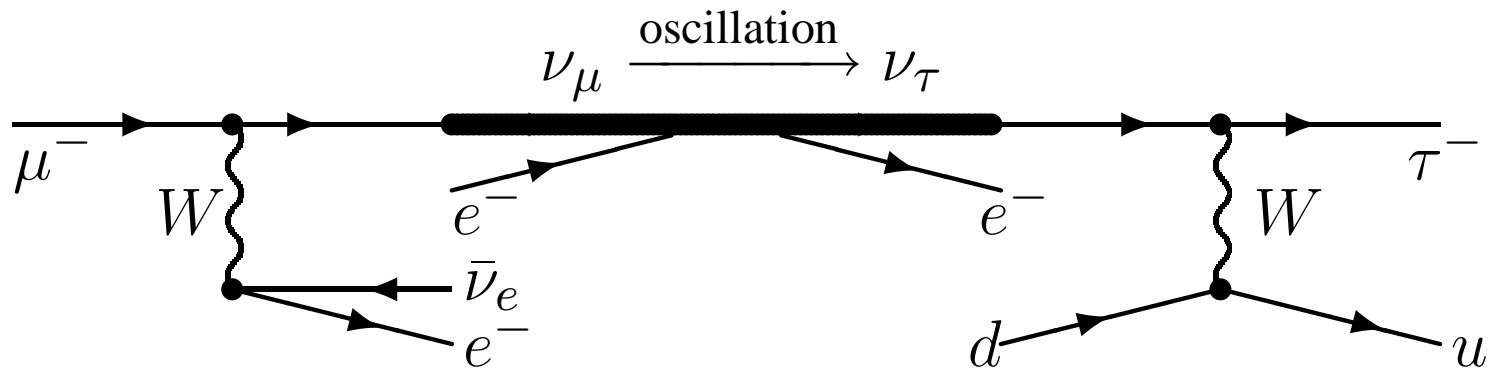


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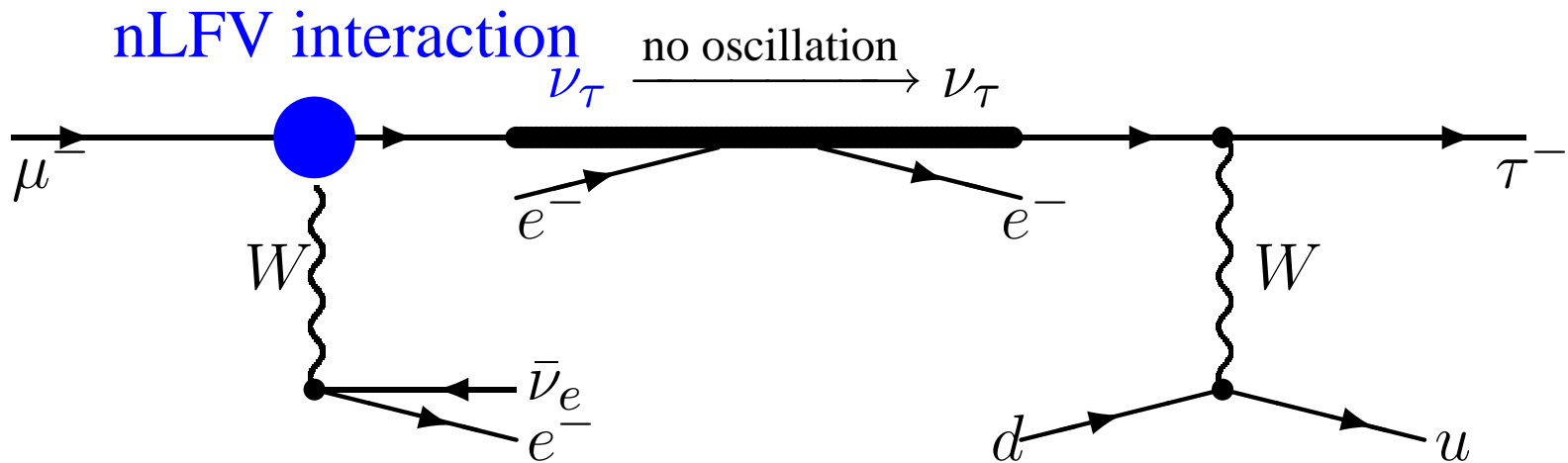
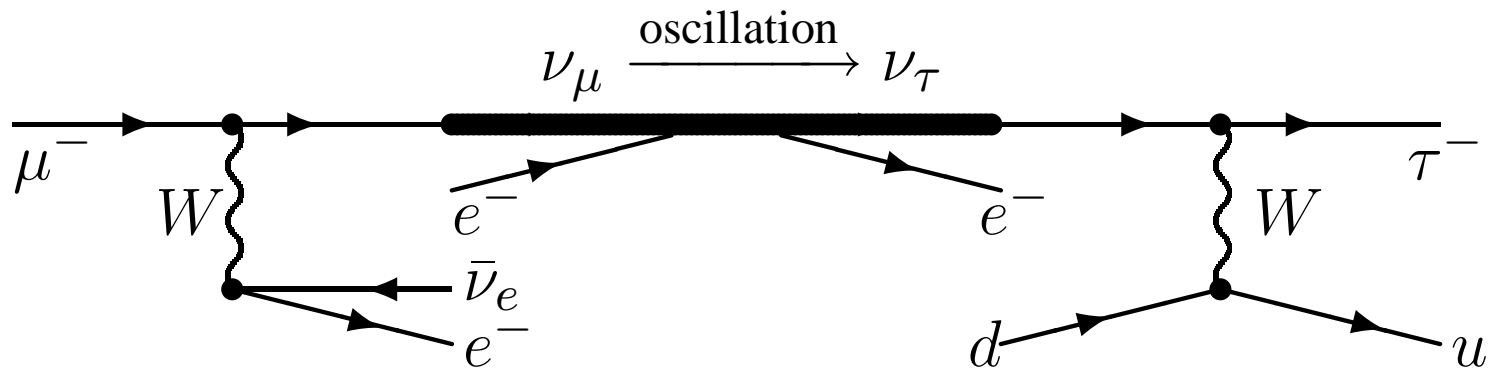
$$\left| \mathcal{A}(\mu \rightarrow \nu_\mu \bar{\nu}_e e) \times \mathcal{A}(\nu_\mu \xrightarrow{\text{osc}} \nu_\tau) \times \mathcal{A}(\nu_\tau d \rightarrow \tau u) \right|^2$$

$$\Gamma_{\text{SM}}(\mu \rightarrow \nu_\mu \bar{\nu}_e e) \times P_{\nu_\mu \rightarrow \nu_\tau} \times \sigma_{\text{SM}}(\nu_\tau d \rightarrow \tau u)$$

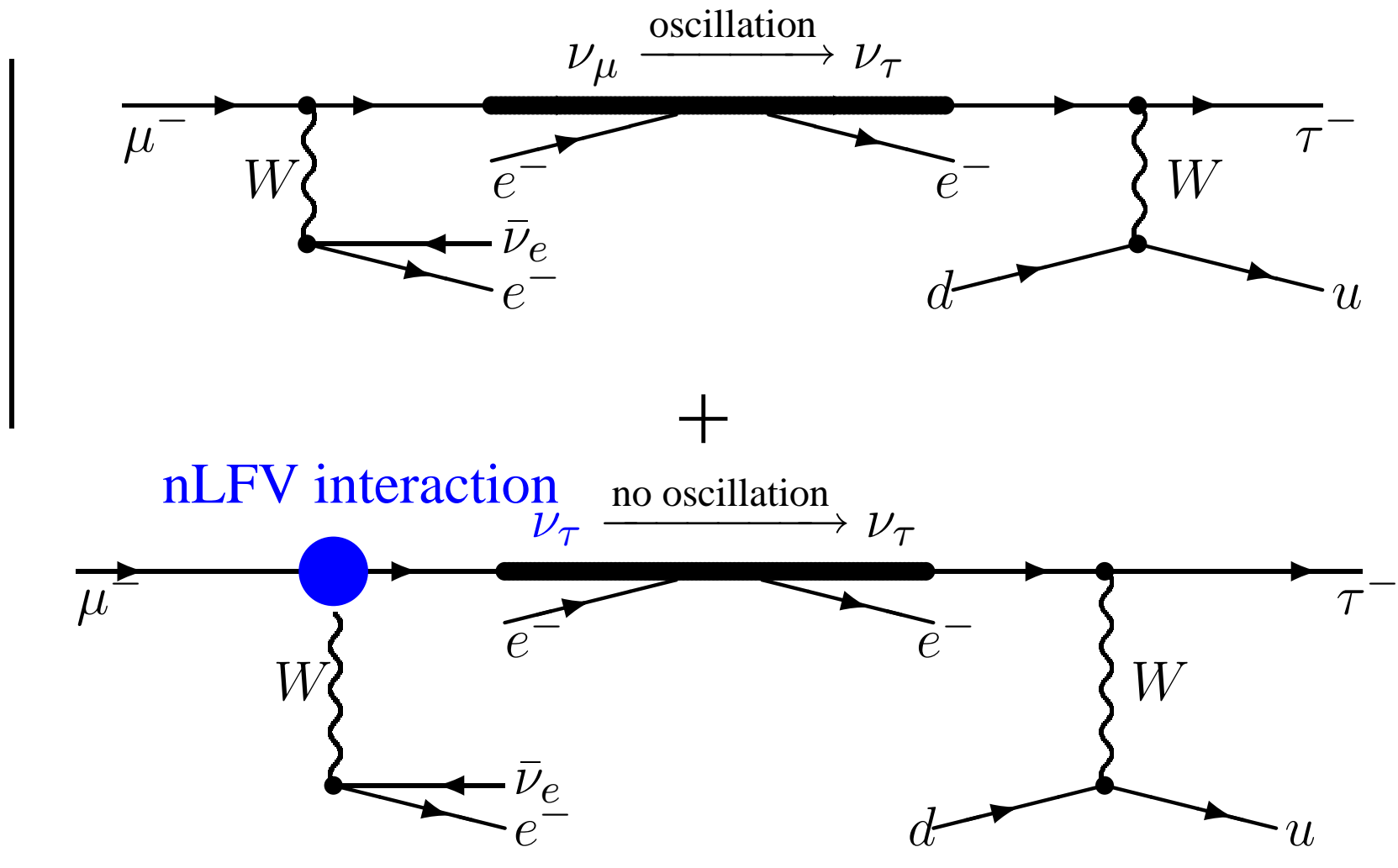
Oscillation with nLFV interactions



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Oscillation with nLFV interactions

$$\left| \begin{array}{c} \nu_\mu \xrightarrow{\text{osc}} \nu_\tau \\ \text{SM} \end{array} \right|^2 + 2\text{Re} \left[\left(\begin{array}{c} \nu_\mu \xrightarrow{\text{osc}} \nu_\tau \\ \text{SM} \end{array} \right) \left(\begin{array}{c} \text{nLFV} \nu_\tau \xrightarrow{\text{no osc}} \nu_\tau \\ \text{SM} \end{array} \right)^* \right]$$

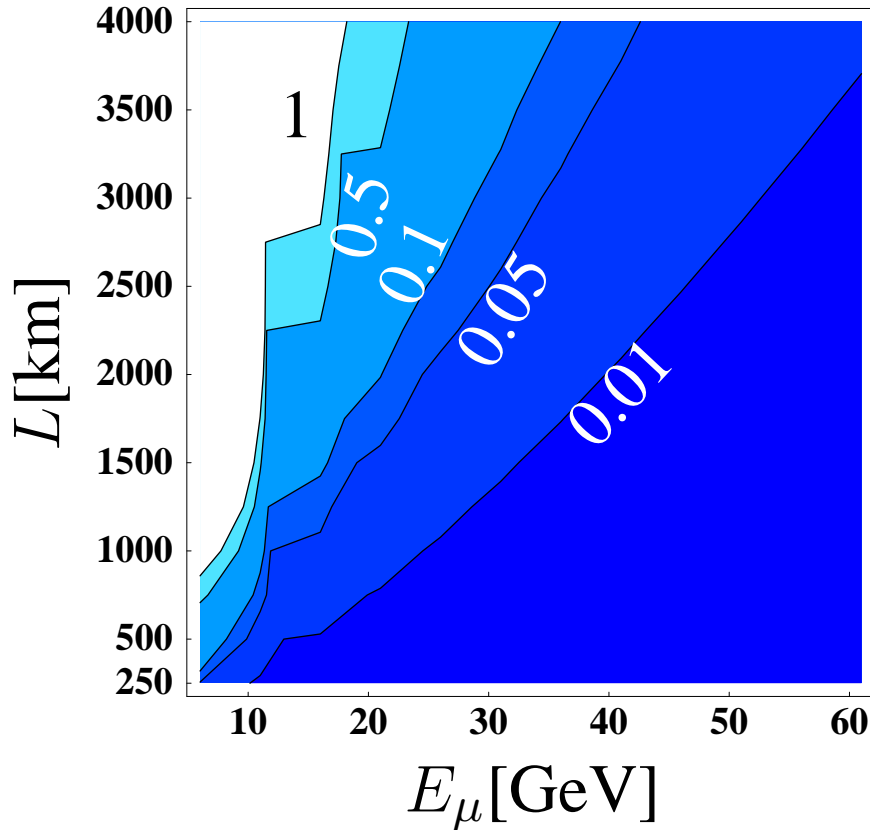
$$= \Gamma_{\text{SM}} \times \left(P_{\nu_\mu \rightarrow \nu_\tau} + 2\text{Re} \left[\epsilon_{\mu\tau}^s \mathcal{A}^*(\nu_\mu \xrightarrow{\text{osc}} \nu_\tau) \mathcal{A}(\nu_\tau \xrightarrow{\text{no osc}} \nu_\tau) \right] \right) \times \sigma_{\text{SM}},$$

where $\epsilon_{\mu\tau}^s \equiv \frac{\mathcal{A}(\mu \rightarrow \nu_\tau \bar{\nu}_e e)}{\mathcal{A}_{\text{SM}}}$.

- The oscillation probability is modified by the interference term due to the nLFV interaction.
- The size of the interference term is $\mathcal{O}(\epsilon_{\mu\tau}^s)$, not $\mathcal{O}(|\epsilon_{\mu\tau}^s|^2)$.
 - This interference effect can occur only in the oscillation process.

Search for the effect of $\epsilon_{\mu\tau}^s$ in $\nu_\mu \rightarrow \nu_\tau$

TO JS Yamashita (2002)



- Necessary (muon) \times (detector size) for 90% CL detection of nLFV.

normalized at $10^{21} \times 100$ [kton]

- $\nu_\mu \rightarrow \nu_\tau$ channel
- $\epsilon_{\mu\tau}^s = 3 \times 10^{-3} e^{i\frac{\pi}{2}}$
- 10% ambiguity is considered in the oscillation parameters, Δm^2 s and θ s, and the CP-phase δ is treated as a free parameter.

$$\mu^- \xrightarrow{G_F} \nu_\mu \xrightarrow{\text{osc.}} \nu_\tau \rightarrow \tau^-$$

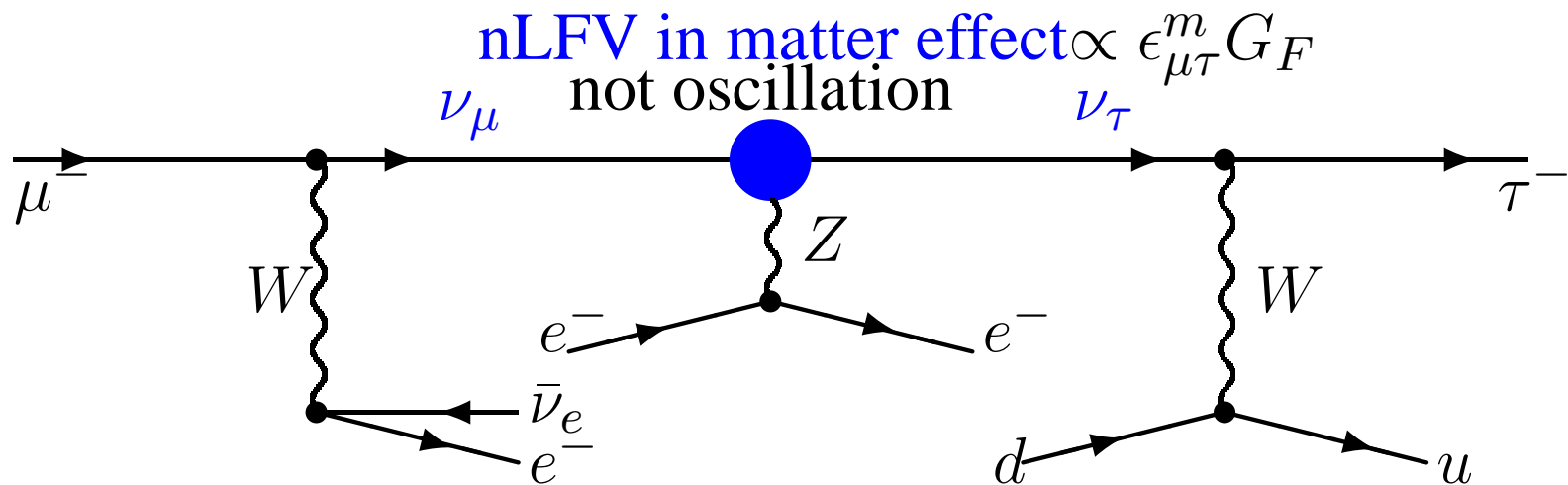
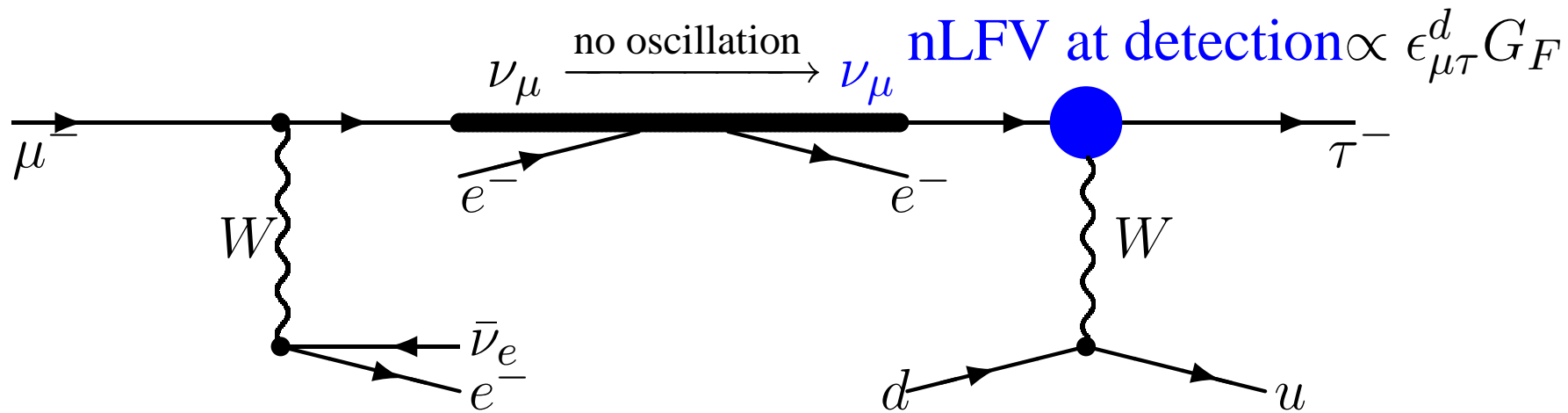
$$\mu^- \xrightarrow{\epsilon_{\mu\tau}^s G_F} \nu_\tau \xrightarrow{\text{no osc.}} \nu_\tau \rightarrow \tau^-$$

- The energy dependence of this signal is $1/E_\nu$.

It is quite different from the standard oscillation effect ($\propto 1/E_\nu^2$).

Oscillation with nLFV interactions

- The other diagrams which contribute to $\mathcal{A}(\mu + I \rightarrow \tau + F)$



Summary of the model independent analysis

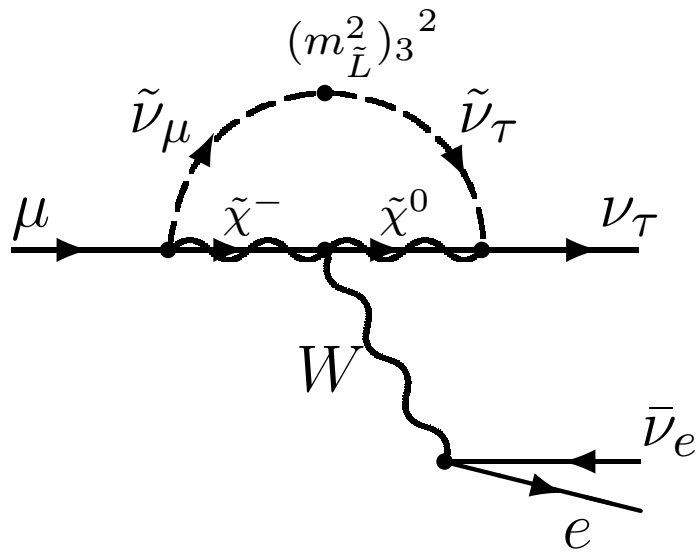
- Detectable signal ...
 - In $\nu_\alpha \rightarrow \nu_\beta$ channel, we can extract the nLFV signals only with $\epsilon_{\alpha\beta}^{s,m,d}$.
 - by using its characteristic energy dependence.
 - the nLFV amplitude *does not include neutrino oscillation* but its final states are the same as the standard one.
 - At a neutrino factory (10^{21} muons \times 100 kt detector), we have a chance to detect the signal of $|\epsilon_{\alpha\beta}^{s,m,d}| \sim \mathcal{O}(10^{-4})$.

Summary of the model independent analysis

- Detectable signal ...
 - In $\nu_\alpha \rightarrow \nu_\beta$ channel, we can extract the nLFV signals only with $\epsilon_{\alpha\beta}^{s,m,d}$.
 - by using its characteristic energy dependence.
 - the nLFV amplitude *does not include neutrino oscillation* but its final states are the same as the standard one.
 - We here deal with $\epsilon_{\mu\tau}^{s,m,d}$ in the $\nu_\mu \rightarrow \nu_\tau$ channel.
- At a neutrino factory (10^{21} muons \times 100 kt detector), we have a chance to detect the signal of $|\epsilon_{\alpha\beta}^{s,m,d}| \sim \mathcal{O}(10^{-4})$.
- We make a numerical calculation for the size of nLFV couplings $\epsilon_{\mu\tau}^{s,m,d}$ in the MSSM with ν_{RS} .

In the MSSM with ν_R s

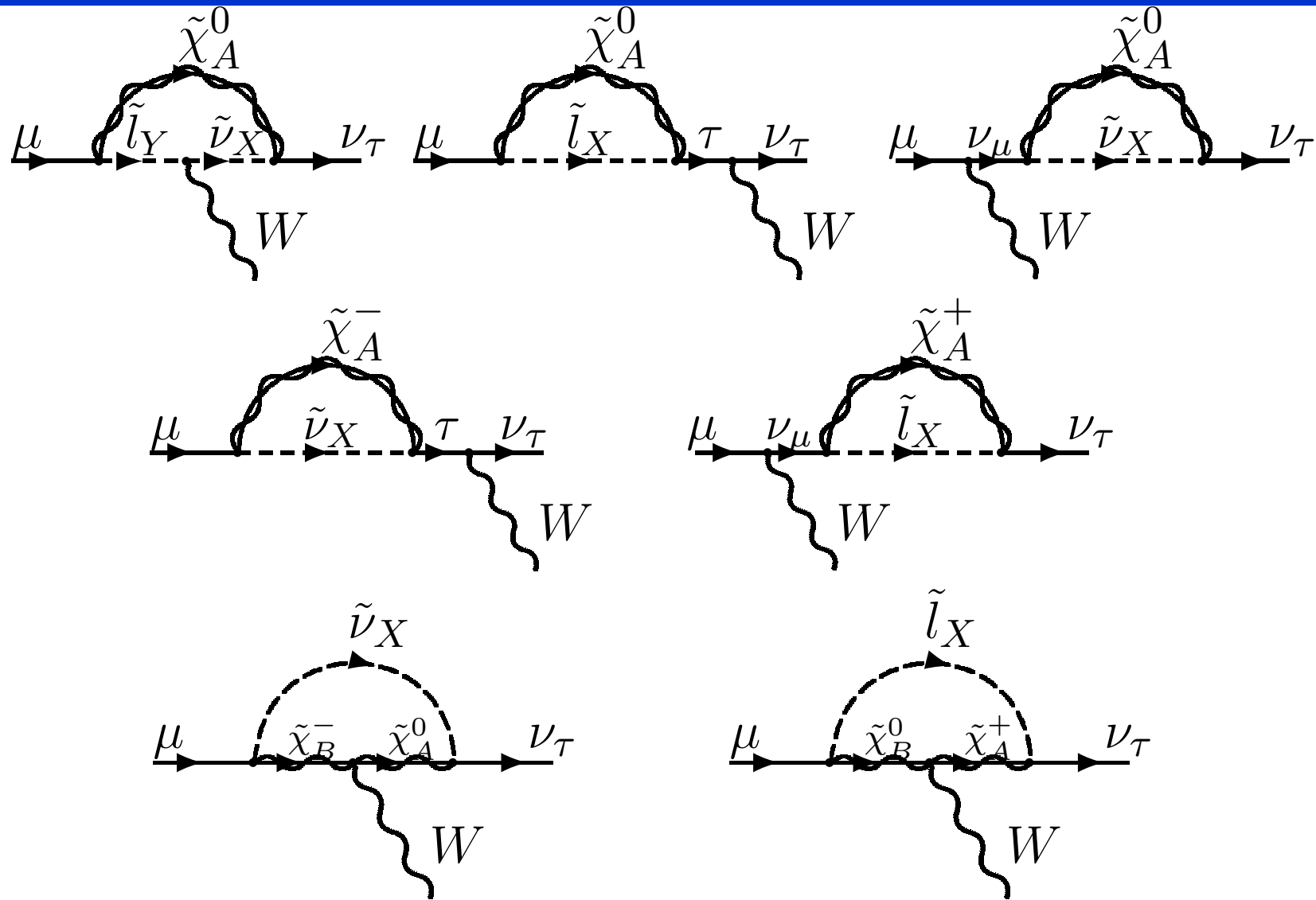
Estimation of $\epsilon_{\mu\tau}^s$ in MSSM + ν_R



$$\sim \frac{g^2}{(4\pi)^2} \frac{(m_{\tilde{L}}^2)_3^2}{m_{\text{SUSY}}^2} G_F \sim \epsilon_{\mu\tau}^s G_F$$

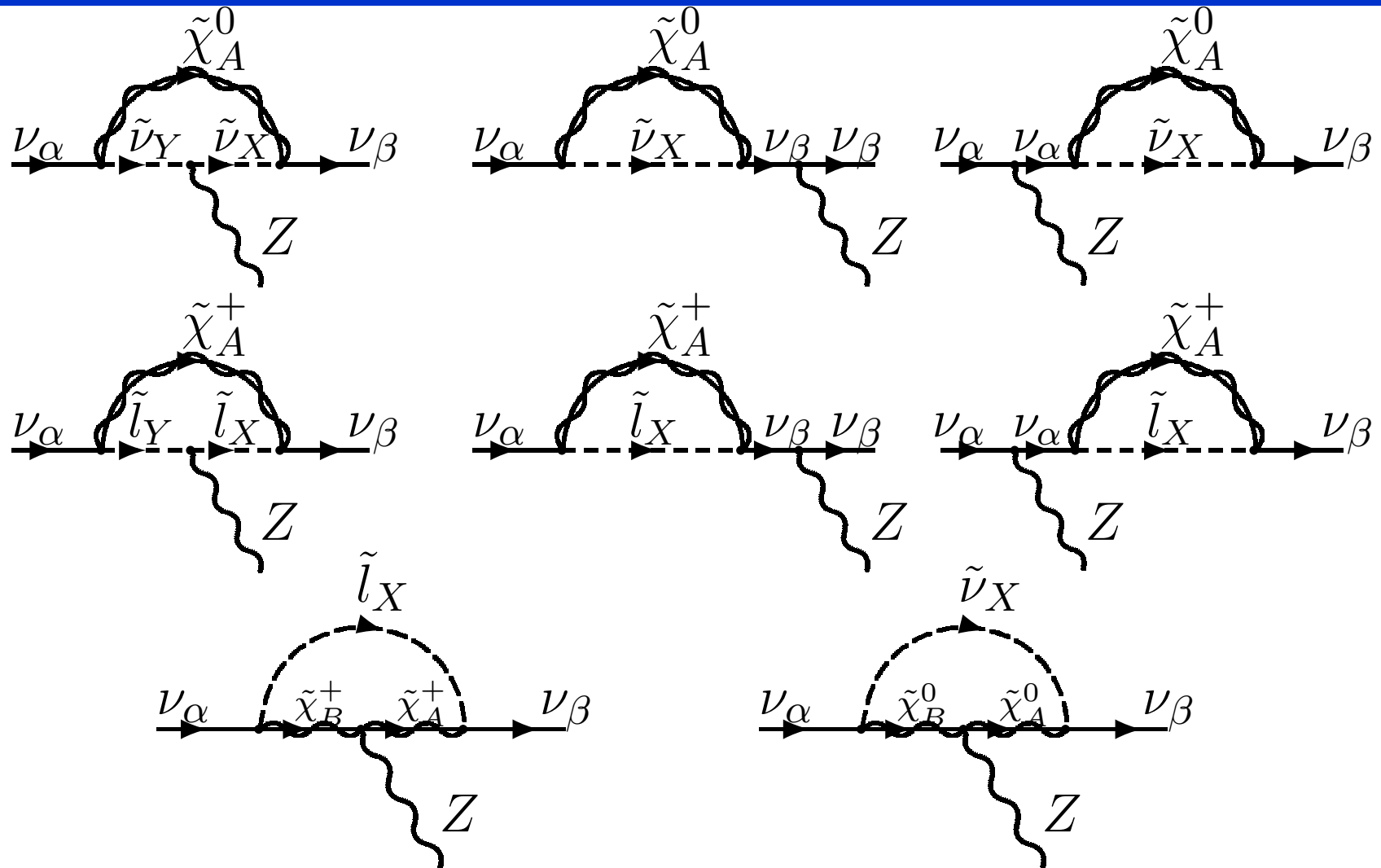
- The origin of the nLFV is the same as that of the cLFV: slepton mixing
- Naively, $\epsilon_{\mu\tau}^s \sim \mathcal{O}(10^{-4})$ in nLFV corresponds to $\text{Br}(\tau \rightarrow \mu\gamma) \sim \mathcal{O}(10^{-8})$ in cLFV.
 - However, for quantitative analysis, it is necessary to make a numerical calculation ...

Numerical evaluation of $\epsilon_{\mu\tau}^s$ in MSSM + ν_R



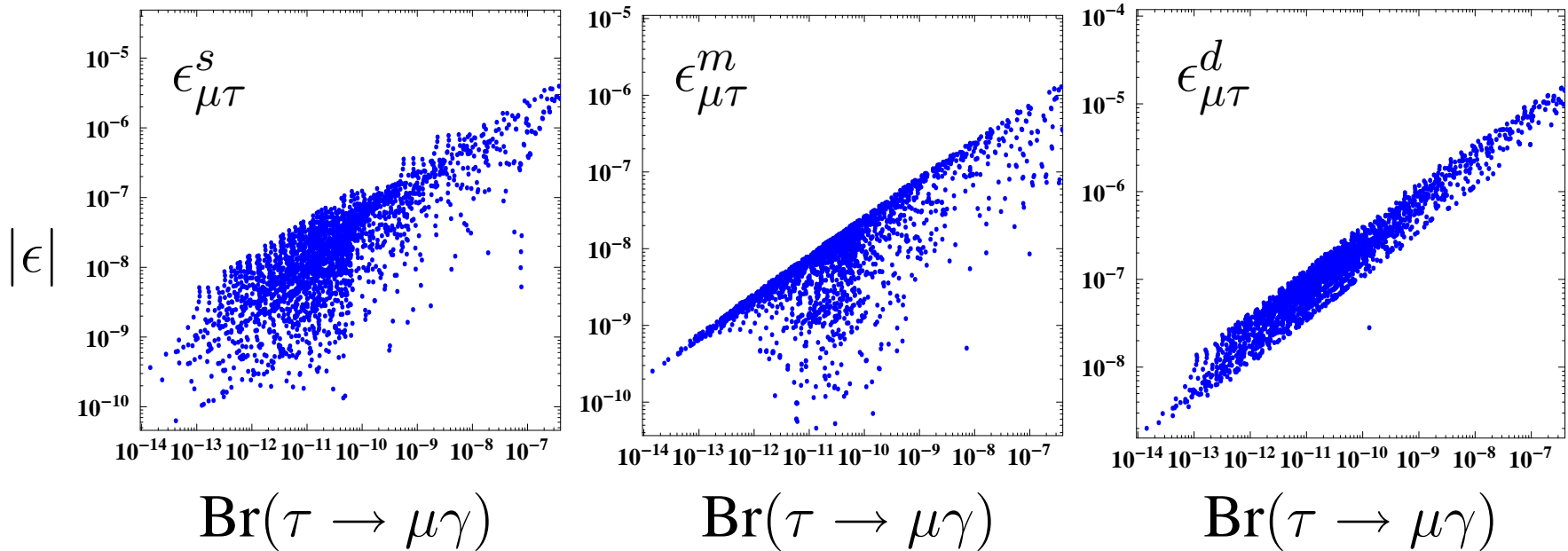
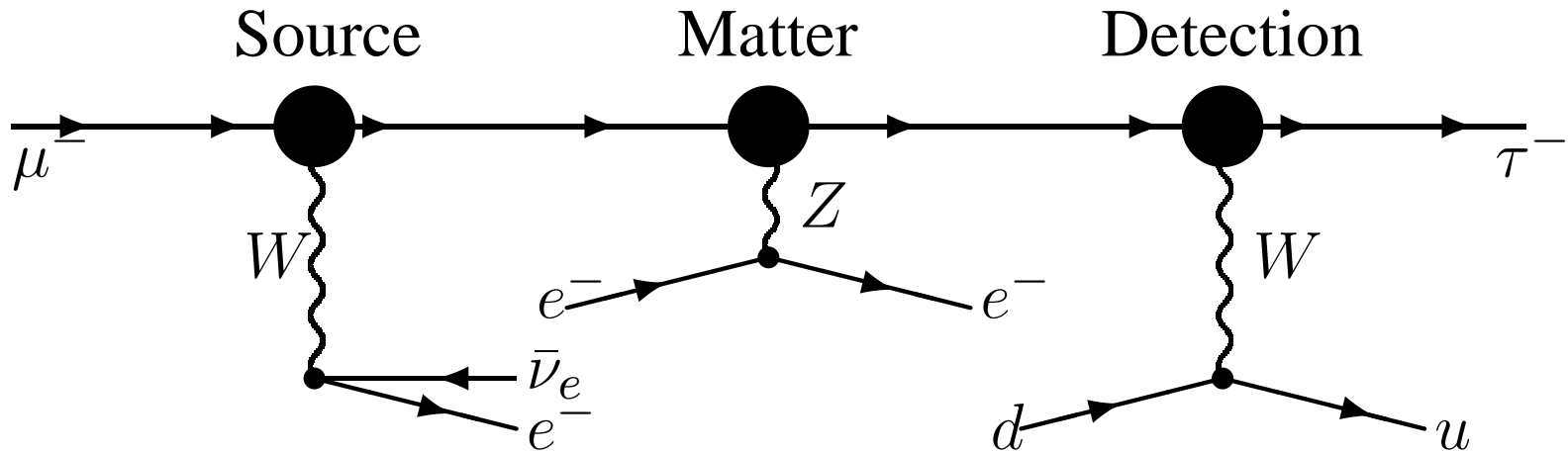
and box diagrams ...

Numerical evaluation of $\epsilon_{\mu\tau}^m$ in MSSM + ν_R

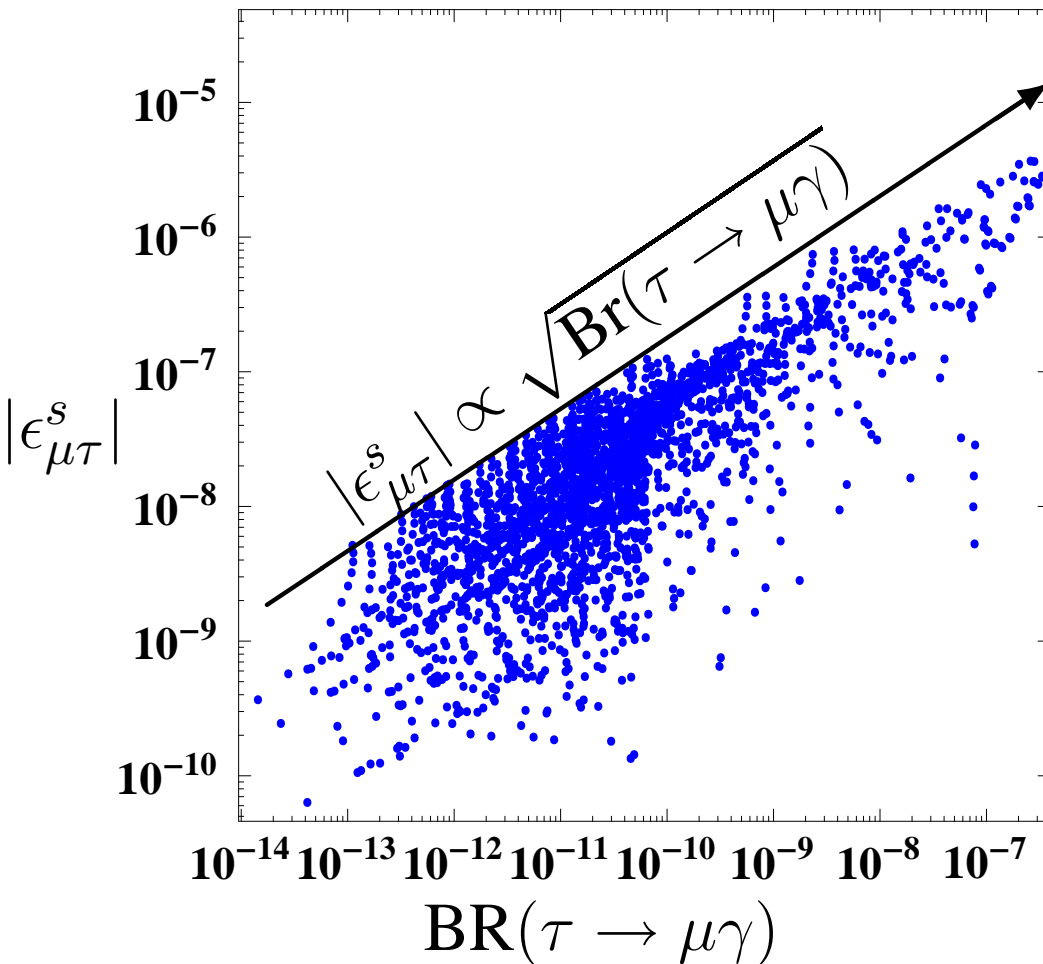


and box diagrams ...

Numerical evaluation of $\epsilon_{\mu\tau}^{s,m,d}$ in MSSM + ν_R



Correlation between nLFV and cLFV



- Correlation between the nLFV coupling $\epsilon_{\mu\tau}^s$ and cLFV process $\tau \rightarrow \mu\gamma$.
 - With some different Y_ν s, we scan the m_0 - $M_{1/2}$ space with $a_0 = 0$, $\tan \beta = 10$, and $\mu > 0$.
- The parameter $\epsilon_{\mu\tau}^s$ is constrained at $\mathcal{O}(10^{-5})$ by the current bound of $\tau \rightarrow \mu\gamma$.
- It is smaller than the naive estimation because of cancellation among diagrams.

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- We evaluate the effective couplings of nLFV which are relevant to $\nu_\mu \rightarrow \nu_\tau$ in the MSSM with ν_R s.
 - $\epsilon_{\mu\tau}^s, \epsilon_{\mu\tau}^m, \epsilon_{\mu\tau}^d$ where $\epsilon_{\mu\tau}^{s,m,d} \equiv \text{exotic/standard}$
 - Detectable size: $|\epsilon_{\mu\tau}^{s,m,d}| \gtrsim \mathcal{O}(10^{-4})$

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- In the MSSM with ν_R s, these couplings are constrained by the process $\tau \rightarrow \mu\gamma$ as $|\epsilon_{\mu\tau}^{s,m,d}| \lesssim \mathcal{O}(10^{-5})$.
- Other model? — R-parity violation ...
 - CPV MSSM ...
 - $SU(5)$ GUT ...